Ultrathin Oxide Films by Atomic Layer Deposition on Graphene

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ABSTRACT: In this paper, a method is presented to create and characterize mechanically robust, free-standing, ultrathin, oxide films with controlled, nanometer-scale thickness using atomic layer deposition (ALD) on graphene. Aluminum oxide films were deposited onto suspended graphene membranes using ALD. Subsequent etching of the graphene left pure aluminum oxide films only a few atoms in thickness. A pressurized blister test was used to determine that these ultrathin films have a Young’s modulus of 154 ± 13 GPa. This Young’s modulus is comparable to much thicker alumina ALD films. This behavior indicates that these ultrathin two-dimensional films have excellent mechanical integrity. The films are also impermeable to standard gases suggesting they are pinhole-free. These continuous ultrathin films are expected to enable new applications in fields such as thin film coatings, membranes, and flexible electronics.

KEYWORDS: Atomic layer deposition, graphene, nanomechanics, thin films

Two-dimensional (2D) materials are promising nanomechanical structures.¹² Graphene, the best known and studied of this class of materials, boasts a high Young’s modulus, intrinsic strength, gas impermeability, and excellent thermal and electrical conductivity.³−⁷ There are numerous applications where flexible ultrathin insulating or oxide films are needed with comparable mechanical properties. The integration of graphene with other 2D or quasi-2D materials may also lead to new functional properties for the composite materials.⁸−¹³ Currently, the range of ultrathin materials is severely limited by the materials and length-scales that are accessible through thin film fabrication.

Mechanical and chemical exfoliation, as well as growth techniques such as chemical vapor deposition, can produce just a handful of ultrathin layered materials.¹⁴−¹⁷ As traditional materials approach ~1 nm film thicknesses, fabrication of freely suspended films is difficult due to stresses or significant voids in the films that destroy the mechanical integrity of the film. To overcome these problems, we use suspended graphene membranes as sacrificial supports to grow high quality ALD films and then remove the graphene to leave the ALD thin film. These experiments demonstrate that ALD on graphene offers a route to create free-standing, ultrathin, quasi-2D structures with atomically controlled thickness and mechanical properties comparable to their bulk counterparts.¹⁸−²¹

Atomic layer deposition films are fabricated using a combination of deposition and etching using a suspended graphene support. The graphene provides an atomically smooth growth surface that can easily be etched away. Graphene is mechanically exfoliated over predefined wells as illustrated in Figure 1a. The graphene is then exposed to a trimethylaluminum (TMA) and NO₂ treatment that forms an adhesion layer for ALD nucleation.²²−²⁴ Aluminum oxide ALD is subsequently grown using TMA/H₂O doses (see Supporting Information). An example of such a graphene/ALD composite film after seven cycles of alumina ALD is shown in Figure 1b. High-resolution cross-sectional transmission electron microscopy on such a graphene sample with a TMA/NO₂ adhesion layer followed by five cycles of TMA/H₂O shows the ALD film to be amorphous and 2.8 ± 0.3 nm thick (see Supporting Information). We then use oxidative etching of the underlying graphene support to leave only the thin alumina ALD film suspended over the predefined well as displayed in Figure 1c. Oxidative etching is carried out in a 1 in. diameter tube furnace at 600 °C with an O₂ gas flow of ~20−40 ccm for ~10 h. This is sufficient to completely etch away the graphene. After the graphene is etched away, the film is no longer visible in the optical microscope, and Raman spectroscopy that shows no signs of a substantial D, G, and 2D peak in the etched samples is used to confirm the absence of graphene (see Supporting Information).

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A pressure difference is applied to the film using a previously reported method where slow diffusion through the SiO2 substrate over pressurizes the film sealed microchamber. An atomic force microscope (AFM) image of such an over pressurized suspended film in Figure 1c is shown in Figure 1d. The ALD film is bulged upward with a maximum deflection through the center of the film, δ = 261 nm, and a radius, a = 2.76 μm, consistent with the radius of the predefined well. At increasing Δp, the film stretches further as δ increases as characterized in Figure 1e. During AFM imaging, the bulge is stable suggesting a constant pressure difference and no significant leak rate of gas out of the microchamber, similar to previous results on graphene membranes. This behavior implies that the aluminum oxide films are pinhole-free and impermeable to the nitrogen gas used for pressurization. The deformation of the film follows

\[ \Delta p = K(\nu)\frac{\delta^3}{a^4}Et \]  

(1)

where E is Young’s modulus, t is the thickness of the film, and \( K(\nu) \) is a constant that depends on the Poisson’s ratio. For the case of aluminum oxide, \( K(\nu = 0.24) = 3.35 \). Figure 2a shows \( K(\nu)(\delta^3/a^4) \) versus Δp for 18 pure alumina ALD films (graphene etched away) fabricated on an exfoliated graphene flake using seven cycles of alumina ALD. The behavior of each film follows a line as expected from eq 1. The average and standard deviation of the slope of these lines gives \( E = 250 \pm 12 \) GPa·nm. The plot of Et versus number of ALD cycles is shown in Figure 2b. A best fit line of the data gives a slope of \( E_{\text{cycle}} = 16.9 \pm 1.4 \) GPa·nm with an intercept of \( E_{\text{ALD,Al2O3}} = 127.1 \pm 13.1 \) GPa·nm. This corresponds to \( E_{\text{ALD,Al2O3}} = 154 \pm 13 \) GPa assuming a thickness gain per cycle of \( t_{\text{cycle}} = 0.11 \) nm.

A similar measurement was performed for a number of different films formed using 4–15 cycles. The plot of Et versus number of ALD cycles is shown in Figure 2a. A best fit line of the data gives a slope of \( E_{\text{cycle}} = 16.9 \pm 1.4 \) GPa·nm with an intercept of \( E_{\text{ALD,Al2O3}} = 127.1 \pm 13.1 \) GPa·nm. This nonzero intercept likely arises from the Et value of the functionalization layer. This slope corresponds to \( E_{\text{ALD,Al2O3}} = 154 \pm 13 \) GPa assuming an ALD growth rate of 0.11 nm/cycle. This Young’s modulus is comparable to previous measurements on much thicker (tens to hundreds of nm) alumina ALD films that have Young’s moduli of 168–220 GPa. Because the films are freely suspended, a mechanical support does not influence the mechanical properties of the ALD thin films. The high Young’s modulus is remarkable considering our samples are 2–3 orders of magnitude thinner than previously measured ALD films.
The pressure-induced strain in the film can be used to tune the mechanical resonance frequency of the suspended films. Figure 3a demonstrates this behavior for a graphene/ALD composite film fabricated using five cycles of alumina ALD. The mechanical resonance is actuated and detected optically as previously reported. We were unable to measure a resonance frequency for the pure alumina ALD films presumably due to the lack of optical reflectivity from these samples. The frequency first decreases and then increases as the film transitions from a bulged upward to a bulged downward state. At sufficiently large pressures far from the minimum frequency, the frequency scales as $f^3 \propto \Delta p$. The slope shows a dramatic decrease in frequency with the addition of alumina ALD cycles as shown in Figure 3b. This behavior can be explained by the pressure-induced changes in the tension in a stretched circular film according to

$$f^3 = 7 \times 10^{-3} \frac{K(u)E}{a^3 \rho_A^2} \Delta p \text{ (2)}$$

where $\rho_A$ is the mass per unit area. From the slope of the lines in Figure 3b and using $K(u)E$ determined by a pressurized blister test on the composite ALD/graphene film (see Supporting Information), we can determine $\rho_A$ of each suspended film before and after each ALD process. All samples showed an increase in $\rho_A$ with the number of alumina ALD cycles as displayed in Figure 3c. The first three cycles showed a larger increase in $\rho_A$ that may be related to the initial nucleation of alumina ALD. The finite $\rho_A$ before any ALD cycles is attributed to the additional mass from the adhesion layer.

Using the measured $\rho_A$ we can estimate the volume density of the ALD films, $\rho_V$, and the areal mass density of the adhesion layer, $\rho_{A_{ad}}$. Because all the samples have an adhesion layer with an unknown $\rho_{A_{ad}}$, we first determine $\rho_V$ from the slope of the lines in Figure 3c for coatings after the nucleation treatment. This determination yields $\rho_V = 2.3 \pm 0.4 \text{ g/cm}^3$ assuming an ALD growth rate of 0.11 nm/cycle. (The anomalously large value at 12 cycles shown in blue was not used in calculating this average and standard deviation.) We then deduce $\rho_{A_{ad}}$ from the measured $\rho_A$ using $\rho_{A_{ad}} = (\rho_A - \rho_V)N$, where $N$ is the number of alumina ALD cycles. This derivation yields an average value and standard deviation of $\rho_{A_{ad}} = 1.4 \pm 0.3 \times 10^{-7} \text{ g/cm}^2$. We can then determine $\rho_V$ for every ALD film in Figure 3c. This procedure yields $\rho_V = 2.4 \pm 0.7 \text{ g/cm}^3$ as shown in Figure 3d. This density is comparable within experimental error to previous densities measured on thicker alumina ALD thin films of 3.0 g/cm$^3$.33
None of the 178 samples fabricated with less than 4 alumina ALD cycles were impermeable to N$_2$ gas after removal of the graphene as shown in Figure 4a. However, the yield of impermeable films increased with number of ALD cycles and reached 85% for 15 ALD cycles as displayed in Figure 4b. This behavior indicates that increasing the number of ALD cycles reduces pinholes or gas diffusion through the film.

For freely suspended films formed using only five cycles of the TMA/NO$_x$ nucleation treatment, AFM images of the films do not show voids (see Supporting Information). This corroborates our measurements of a contribution from the adhesion layer to $E$ and $\rho_{\text{Alumina}}$ in Figures 2b and 3c. Future work will examine the dependence of the adhesion layer and its role in nucleating continuous pinhole-free ALD alumina thin film growth on graphene.

In conclusion, a new class of ultrathin films has been created based on aluminum oxide ALD on graphene. These films are mechanically robust, pinhole-free, and have approximately nanometer thicknesses while still maintaining a Young’s modulus comparable to their much thicker counterparts. The manufacturability, thickness control, and versatility of the ALD process means that materials and processing can be tailored to suit many applications where traditional silicon or graphene-based thin film mechanical devices fail to offer the needed functionality.$^{14,35}$ Furthermore, these films can be integrated with graphene or other nanomechanical structures to create multifunctional quasi-2D electromechanical structures.

**REFERENCES**


(7) Zhang, Y. B.; Tan, Y. W.; Stormer, H. L.; Kim, P. Nature 2005, 438, 201–204.


**ASSOCIATED CONTENT**

1 Supporting Information

Experimental methods, TEM imaging of a graphene/ALD composite, Raman spectrum verifying the etching of graphene, elastic constants of pure ALD films, initial tension in graphene and graphene/ALD composite films, and pure ALD films from the nucleation treatment. This material is available free of charge via the Internet at http://pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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